

ME 531 April 15 2020

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Monday Sameer is going to present crystal plasticity example.

- (1) Latent Hardening (Theory)
- (2) Geometric hardening,  $k_0, k_1$   
Eqn 3.12.14
- (3) Stress Rate (Objective) WARF3D  
- Jaumann

Rice Asaro Needleman

$$H_{\alpha\beta} = q H + (1-q) H_{S_{\alpha\beta}}$$

$q$  = material constant

$$\dot{\epsilon}^{\alpha} = \sum_{\beta=1}^N H_{\alpha\beta} \dot{\gamma}(\beta)$$

$$H_{\alpha\beta} = \begin{bmatrix} H & H \\ H & H \end{bmatrix}$$

2 slip systems

$$q = 1$$

System 2  
System 1  
 $\delta$   
 $\bar{\tau}$  increasing



$\tau$  increasing (isotropic hardening)  
in 2 same amount  
as 1, H governs

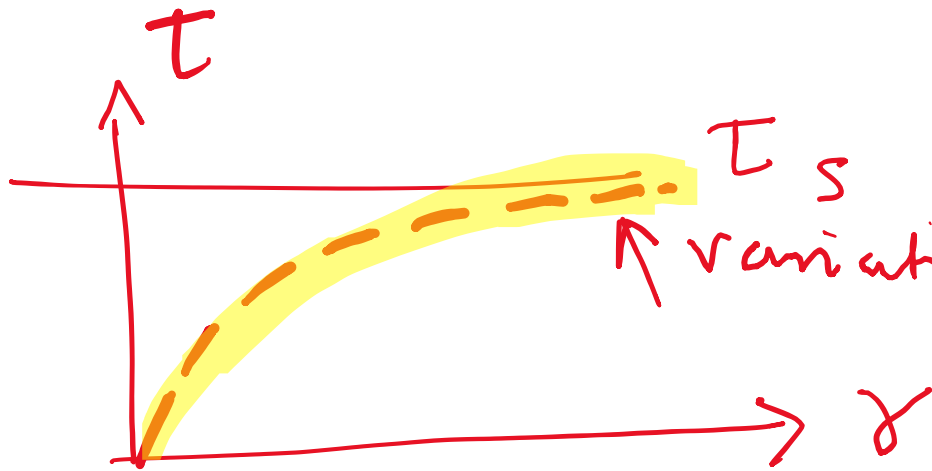
$$H_{\alpha\beta} = q H + (1 - q) H \delta_{\alpha\beta}$$

$q = \text{const. (idealization)}$

$$H(\gamma) = H_0 \operatorname{sech}^2 \left( \frac{H_0 \gamma}{\tau_s - \tau_0} \right)$$

= PAN

= changes with  $\gamma$   
 ( $H_0, \tau_0$  current value of shear resolved  
 $\tau_s = \text{saturation value}$ )



$$\tau = H \gamma \quad \text{single system}$$

$H$  can be determined from experiments on single crystals oriented to activate a single slip system

⇒ Upon integration

$$\tau(\gamma) = \tau_0 + (\tau_s - \tau_0)$$

$$\tanh \left[ \frac{H_0 \gamma}{\tau_s - \tau_0} \right]$$

WARP 3D  $\Rightarrow$  MTS model for  
hardening

$$\tau \approx \tau_a + \tau_i + \tau_\gamma$$

(single system)      ( $\epsilon, T$  effects)  
midterm

Latent Hardening ?

LHR = Latent hardening  
ratio

=  $L_{ji}$   
latent system  $j$ , primary system,  $i$

$$L_{ji} = \tau_j^0 / \tau_i^{\max}$$

$\tau_j^0$  = yield stress of latent system

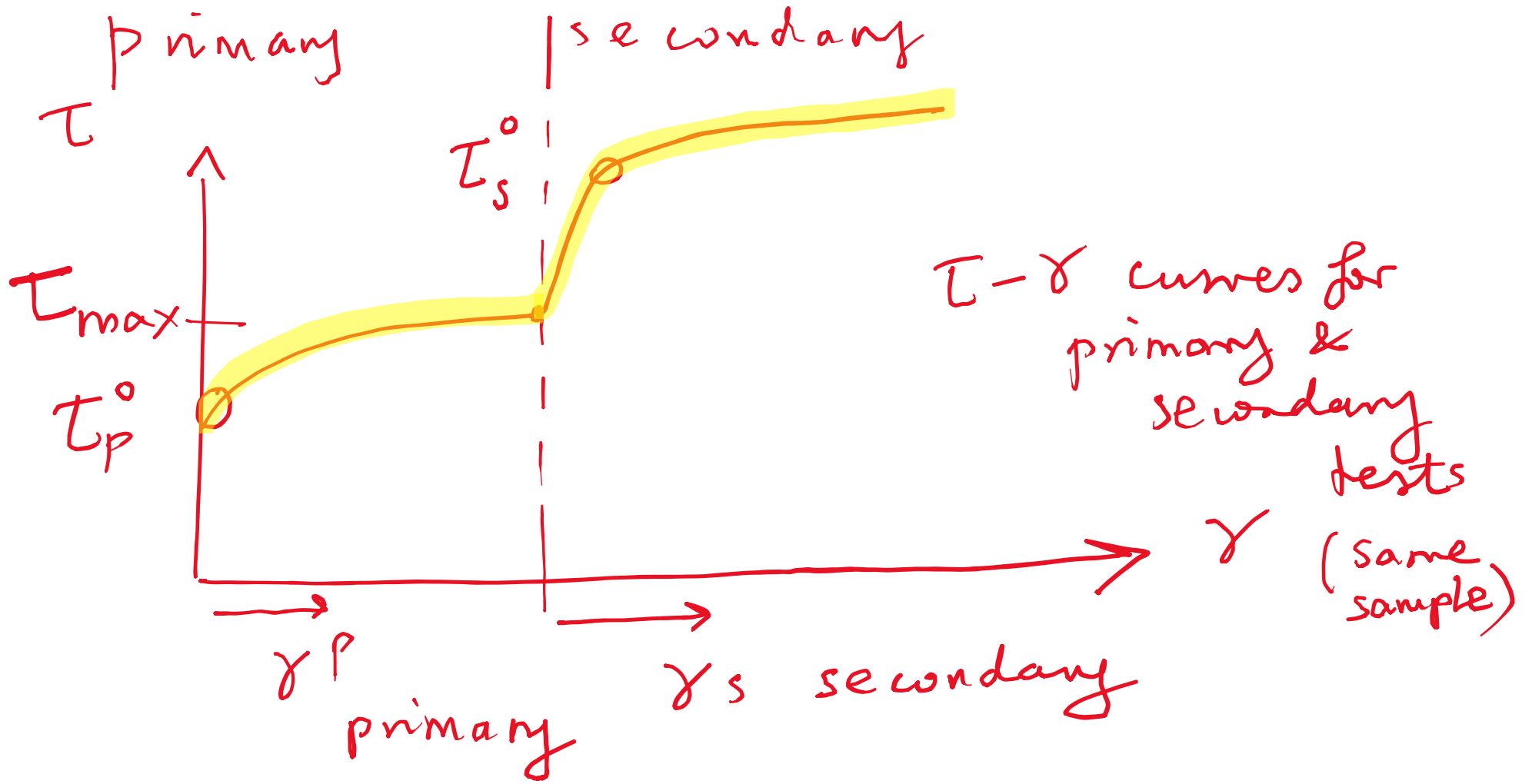
$\tau_i^{\max}$  max stress reached in the primary system

Primary deformation

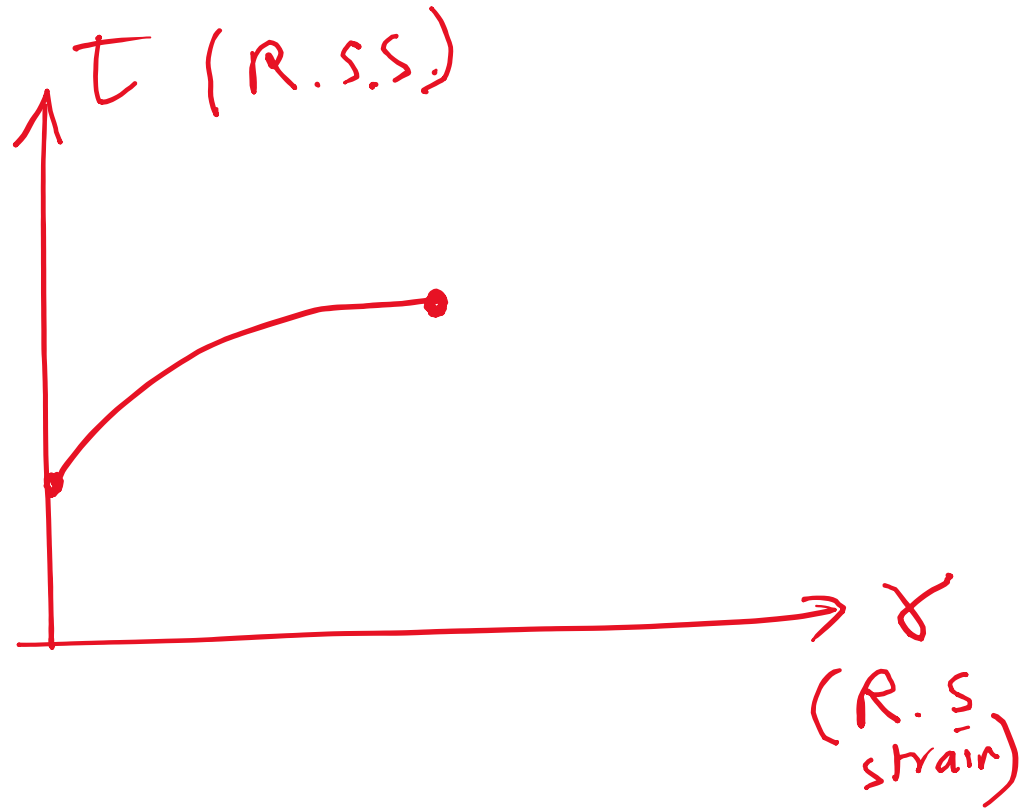
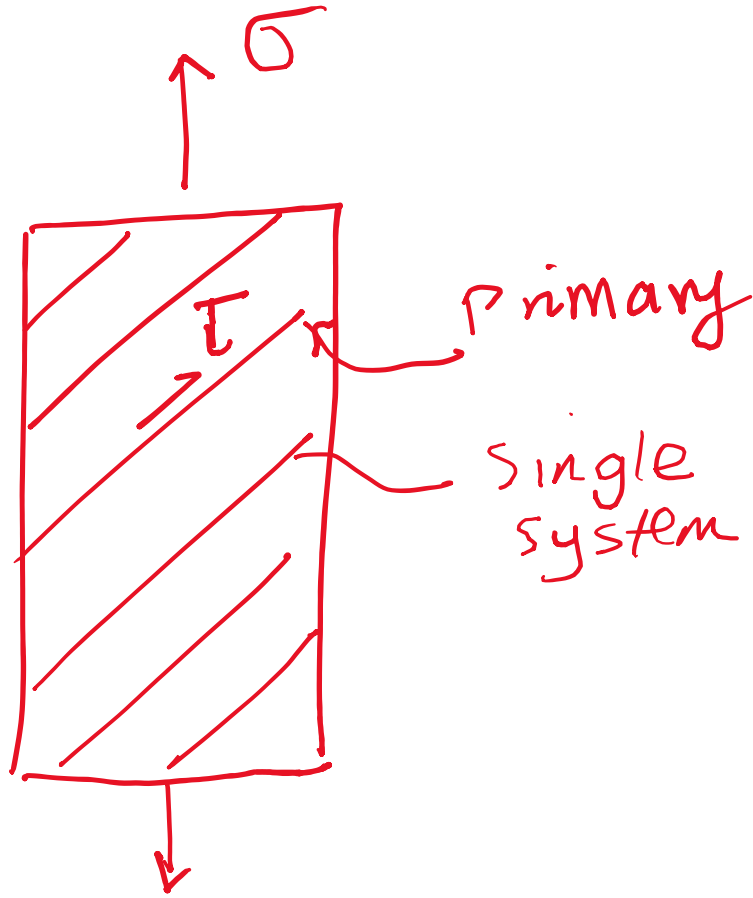
— easy glide — only one system is activated

Secondary test — latent system is active interacts with primary system.

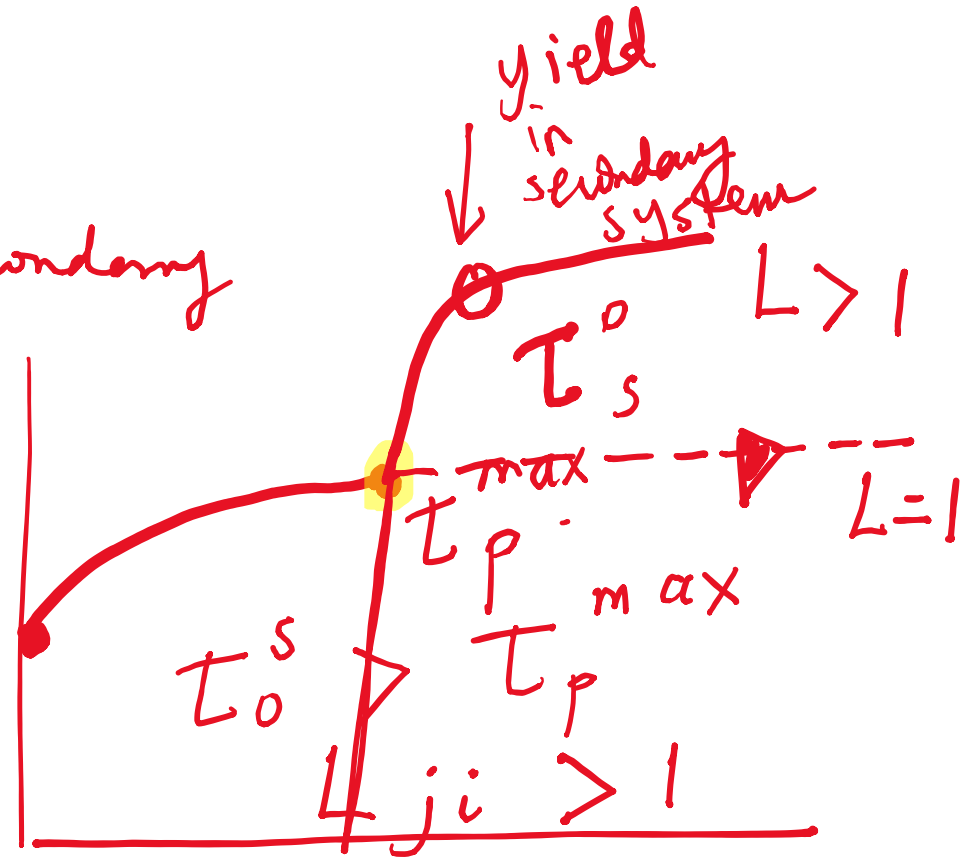
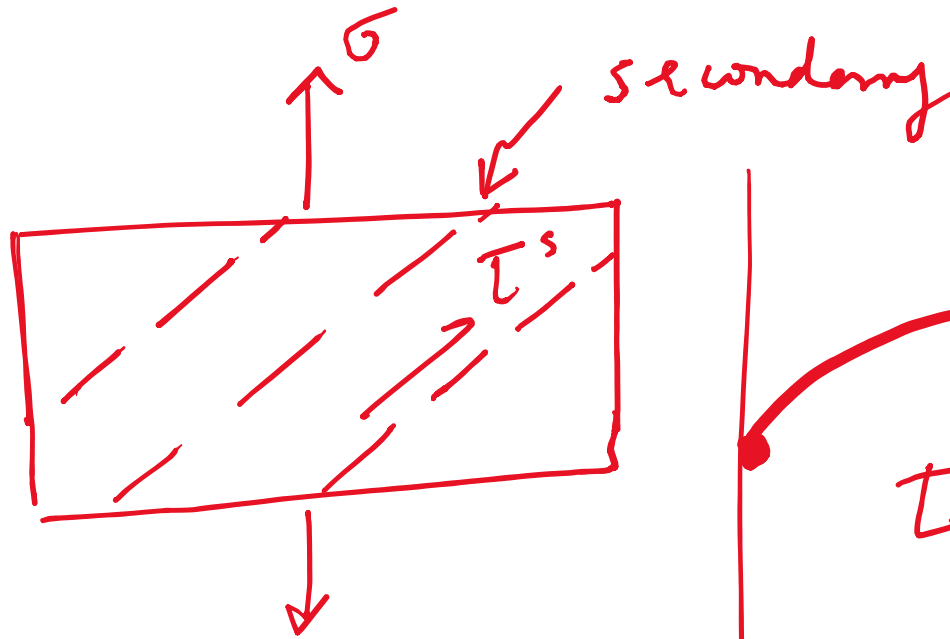




# Primary test



Rotates the sample



L a tent hardening (cross hardening)

In real life

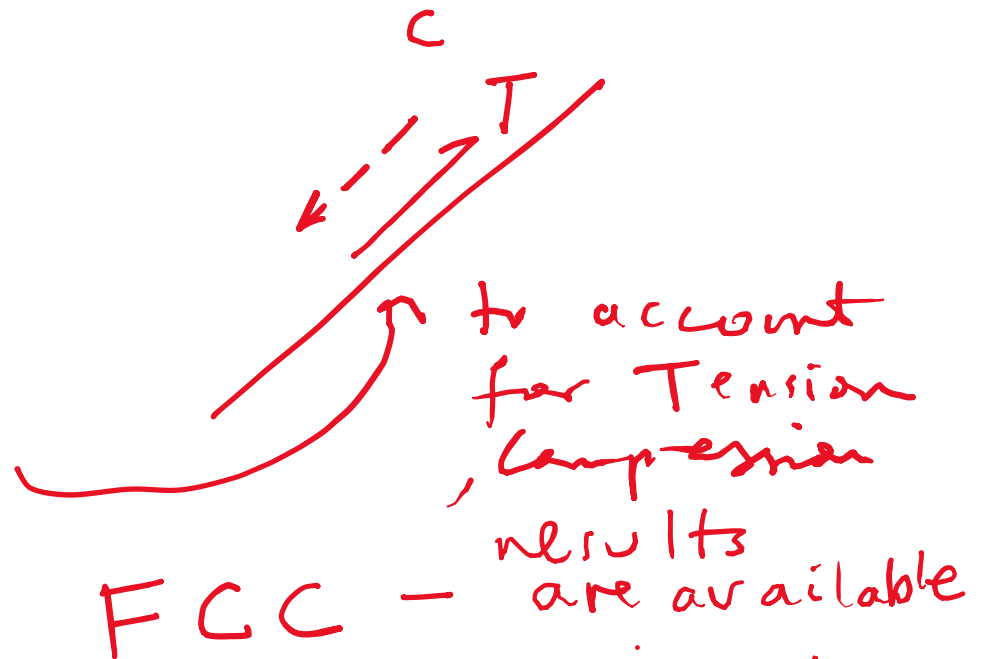
$H_{\alpha\beta}$  12x12 matrix

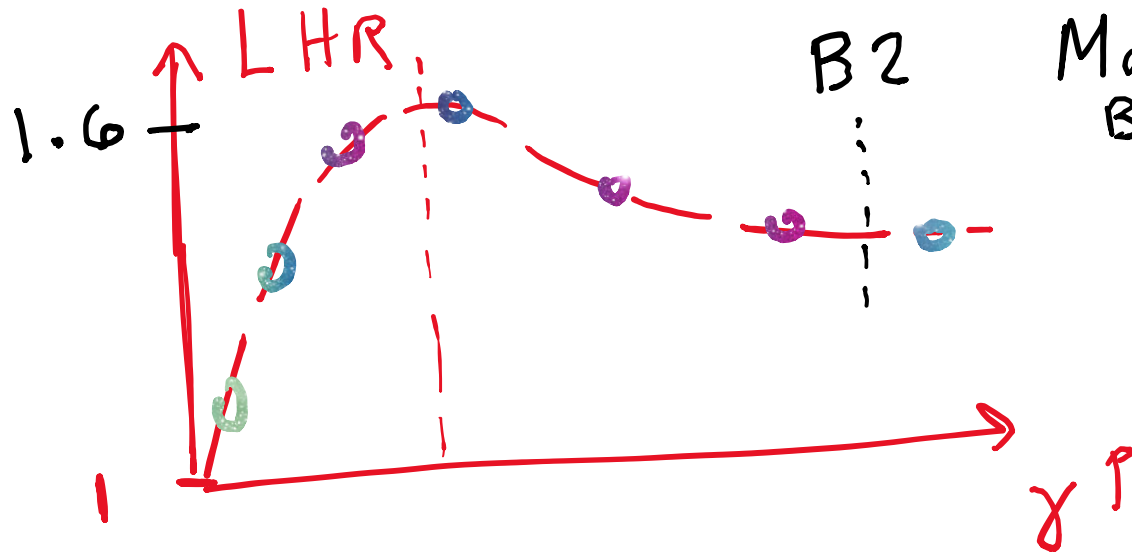
In fact

$H_{\alpha\beta}$  24x24 matrix

(symmetries)

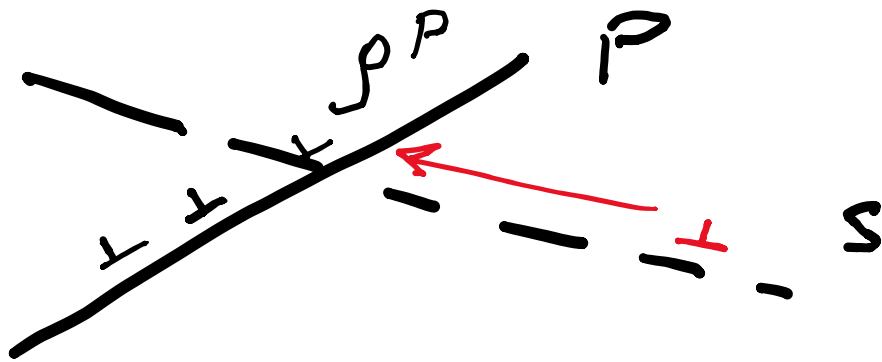
For BCC, HCP it is more complicated





Mand  
Berveiller,  
Francois Zaoui,

FCC ] some results.  
BCC ] totally  
HCP ] not understood



A to D

A  $(\bar{1}11)$

B  $(111)$

C  $(\bar{1}\bar{1}1)$

D  $(1\bar{1}1)$

numbers 1 to 6

$\pm [011]$

$\pm [101]$

$\pm [1\bar{1}0]$

$\pm [0\bar{1}1]$

$\pm [\bar{1}01]$

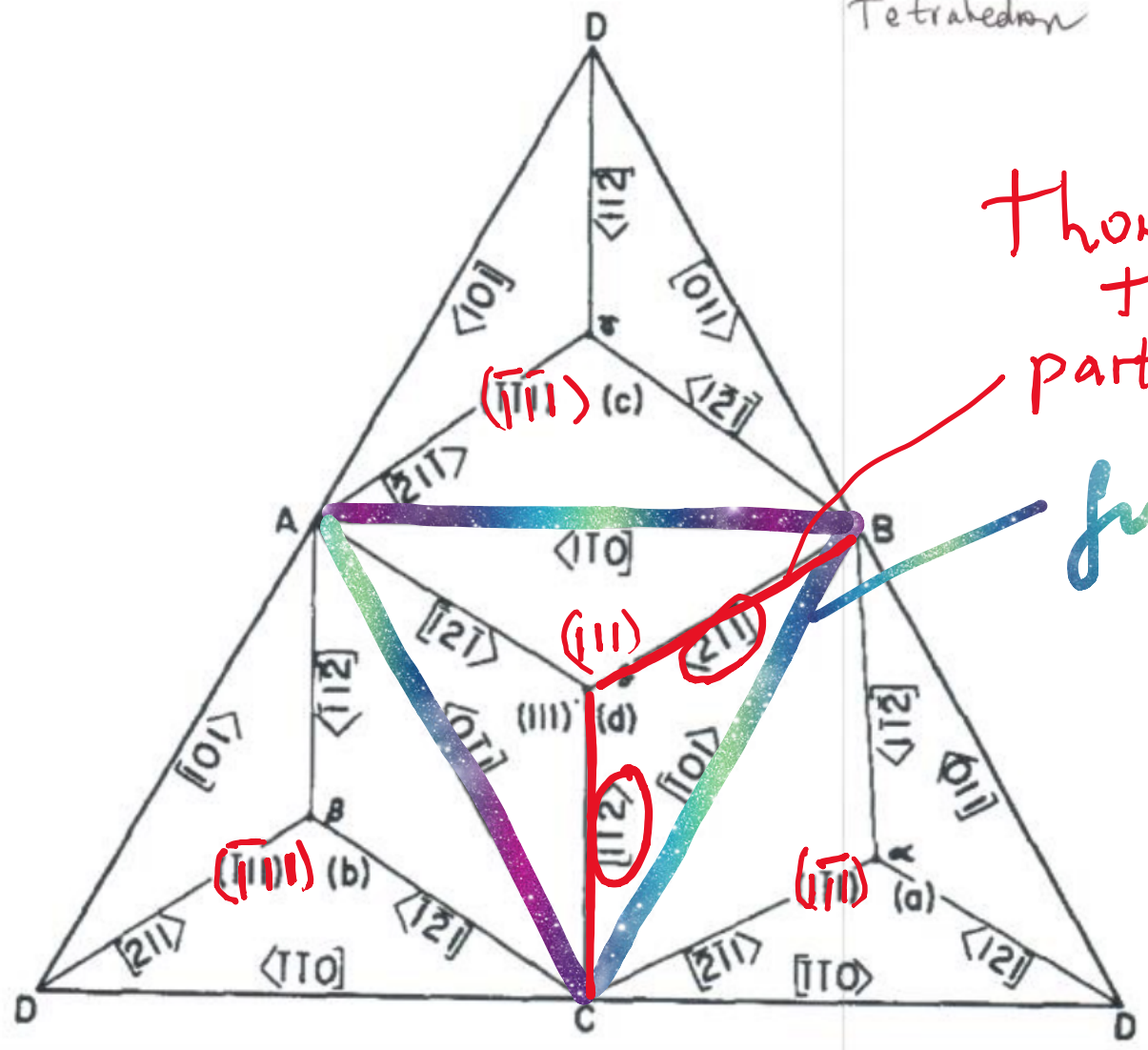
$\pm [110]$

are slip directions

Stereographic projection  
has 24 standard triangles  
each representing one slip  
system.

- 1) Thompson tetrahedra
- 2) Stereographic projection

Thompson Tetrahedron



Thompson tetrahedra partial  
full } FCC



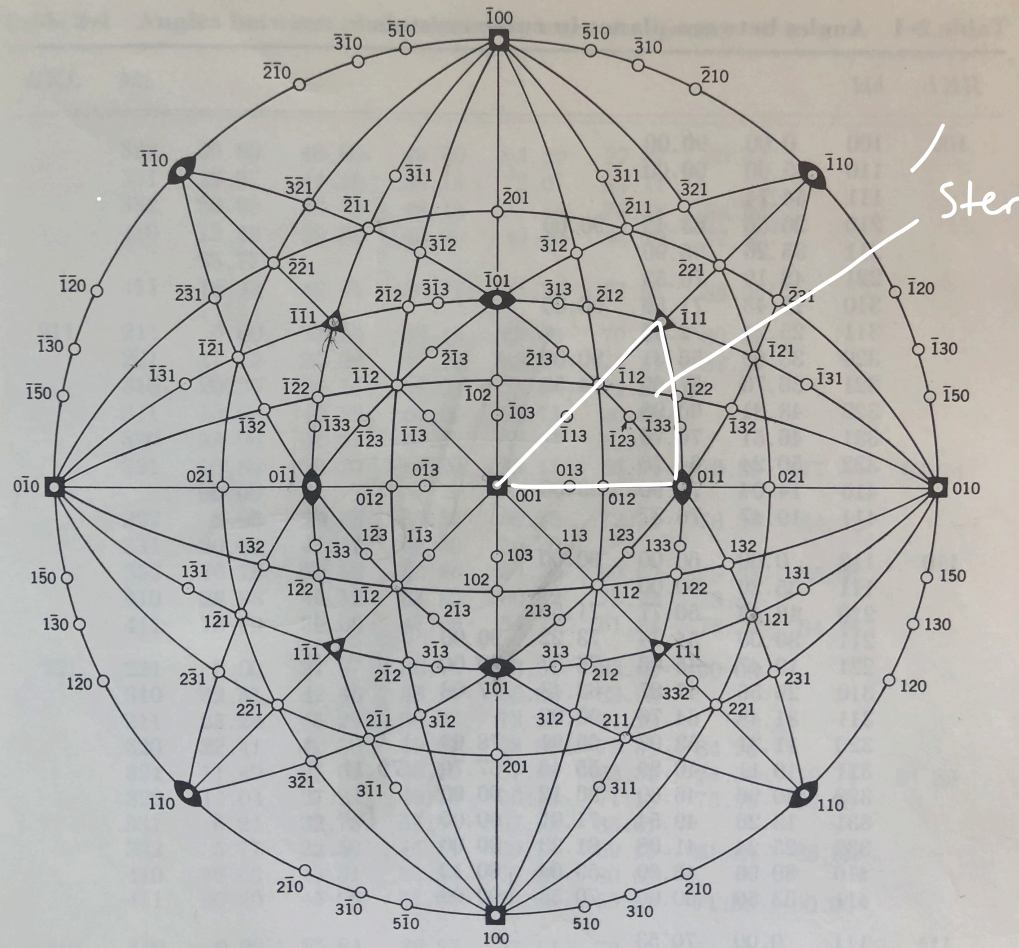
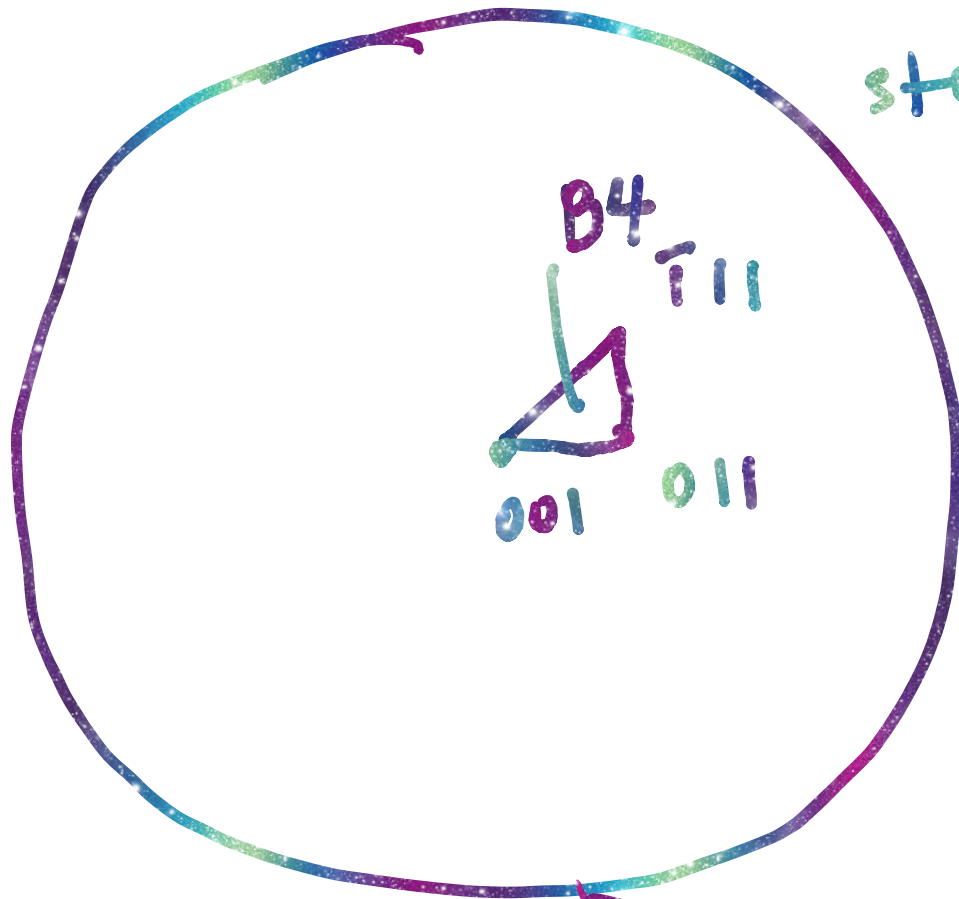


Fig. 2-9 Standard (001) stereographic projection of poles and zone circles for



stereographic triangle

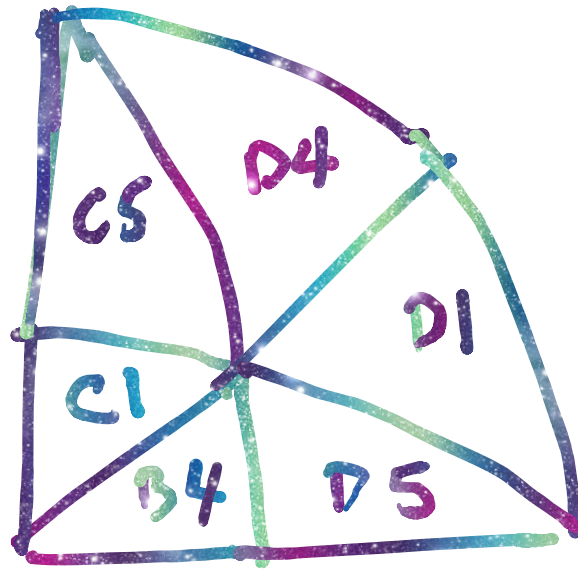
B4 (111) [101]  
 (111) [101]

primary

(self hardening)

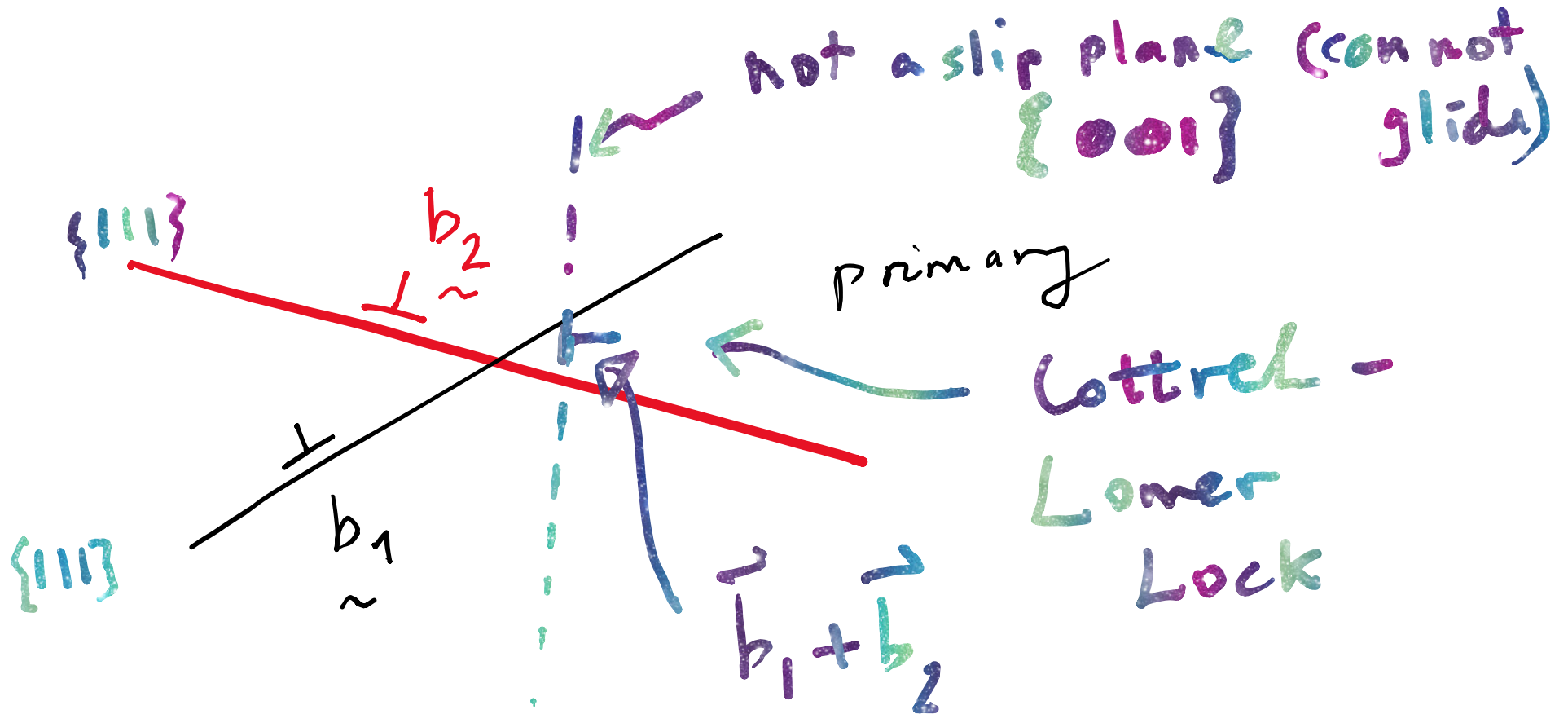
primary  
 (111)

stereographic projection



Dislocations in the primary system  
 can interact with conjugate or  
 critical systems to create  
 Lomer Cottrell dislocations  
 if the resulting dislocation has slip  
 planes that lies on neither one

of the intersecting planes.



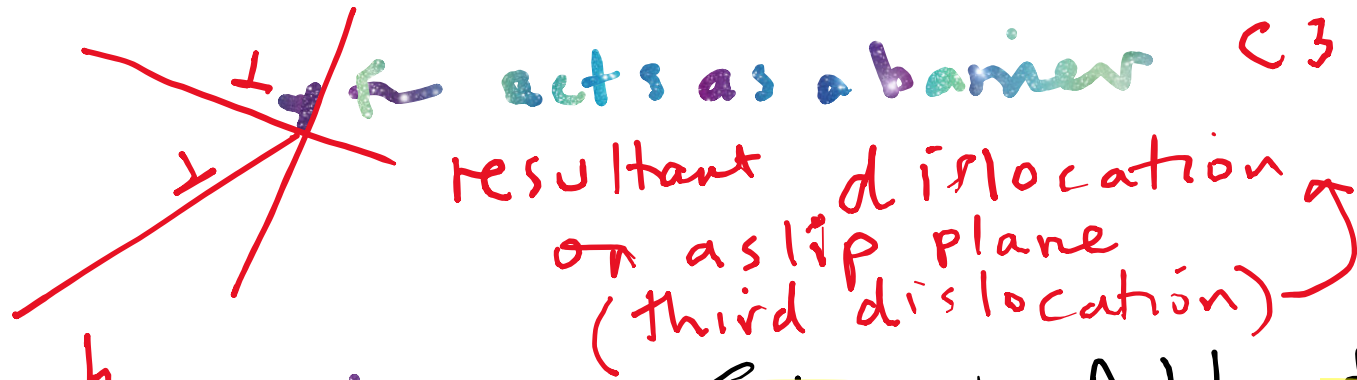
# L - C sessile lock formation

(D 4 primary)

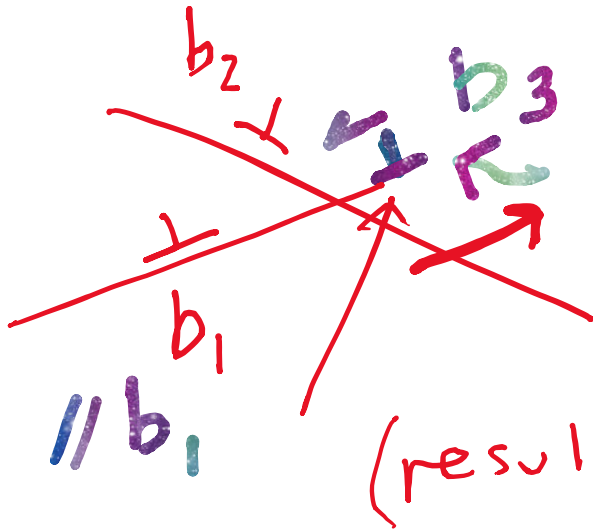
		Angle Between Two planes	resultant ↓ - b	SP slip plane ↓ (100)
1) A6	$(\bar{1}11)[110]$	$70.5^\circ$	$[011]$	
A6	$(\bar{1}11)[\bar{1}10]$	$70.5$		
C1	$(\bar{1}\bar{1}1)[0\bar{1}1]$	$109.5$	$[\bar{1}\bar{1}0]$	$(001)$
C1	$(\bar{1}\bar{1}1)[011]$	$109.5$	$[\bar{1}12]$	$(110)$

1) Lomer Cottrell Lock (strongest)

2) Hirth Lock A3



3) Glissile Attractive Junction  
Sign is such attractive  
(Force needed to separate from the junction)  
(resultant dislocation on either slip plane)



B4 primary +  $\frac{A2, C5, D1, D6}{\text{secondary}}$   $\Rightarrow$  glisite attractive Junct.

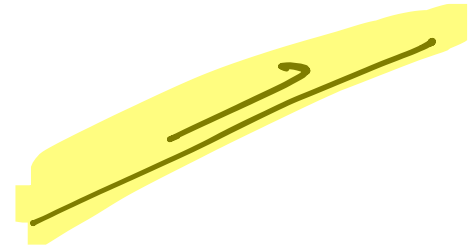
(4) Crosslip in system D4

(5) Coplanar system B2 & B5

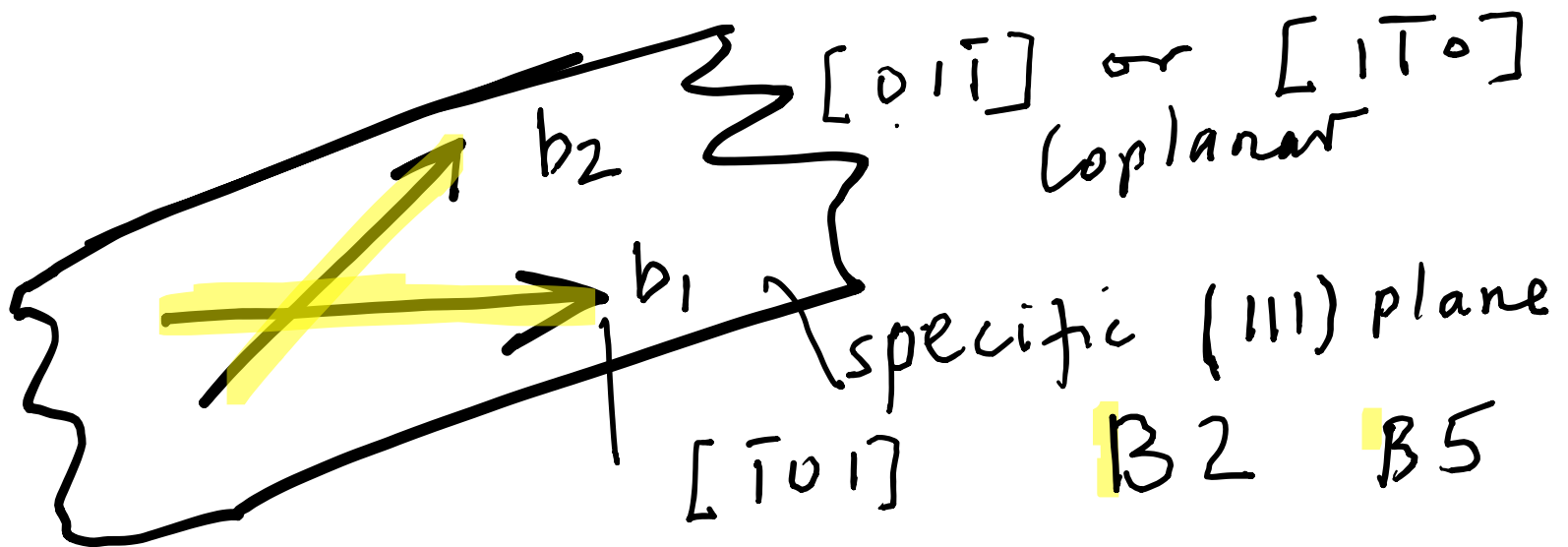
(6) Self hardening in system  
B4

# (6) Self Hardening

B4 (111)  $[\bar{1}01]$  0  
(111)  $[10\bar{1}]$  180



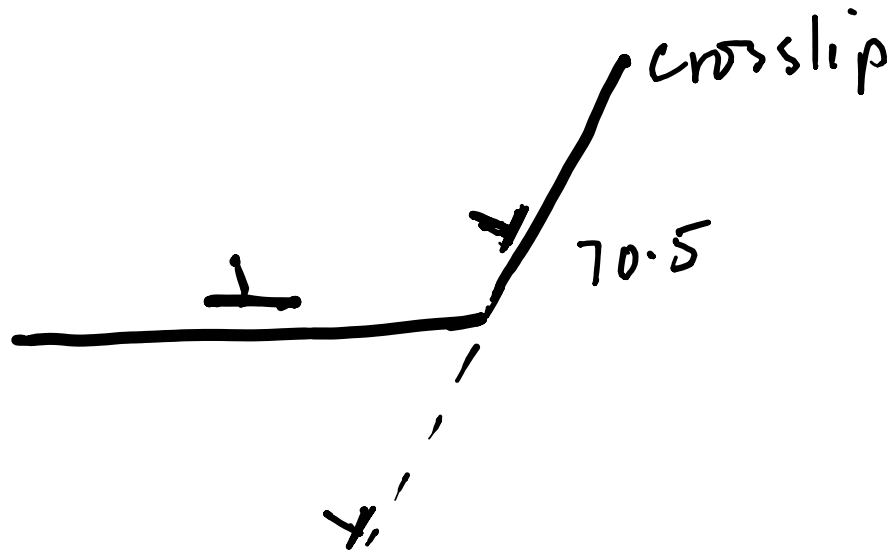
(5)





(4) Cross slip

D4	(1 $\bar{1}$ 1)	[ $\bar{1}$ 01]	70.5
D4	(1 $\bar{1}$ 1)	[10 $\bar{1}$ ]	70.5



## Further Analysis:

$$T_j = \sum \underbrace{h_{ji}} \underbrace{\delta_i}$$

Primary system active step I

$$T_1 = h_{11} \delta_1$$

$$T_2 = h_{21} \delta_1$$

only  
 $\delta_1$  active

$T_1$  primary (P)  $T_2$  secondary (S)

Step II ( $\gamma_2$  active only)

$$\tau_2 = h_{22} \gamma_2 + h_{21} \gamma_1^{\max}$$

$\uparrow$   
active

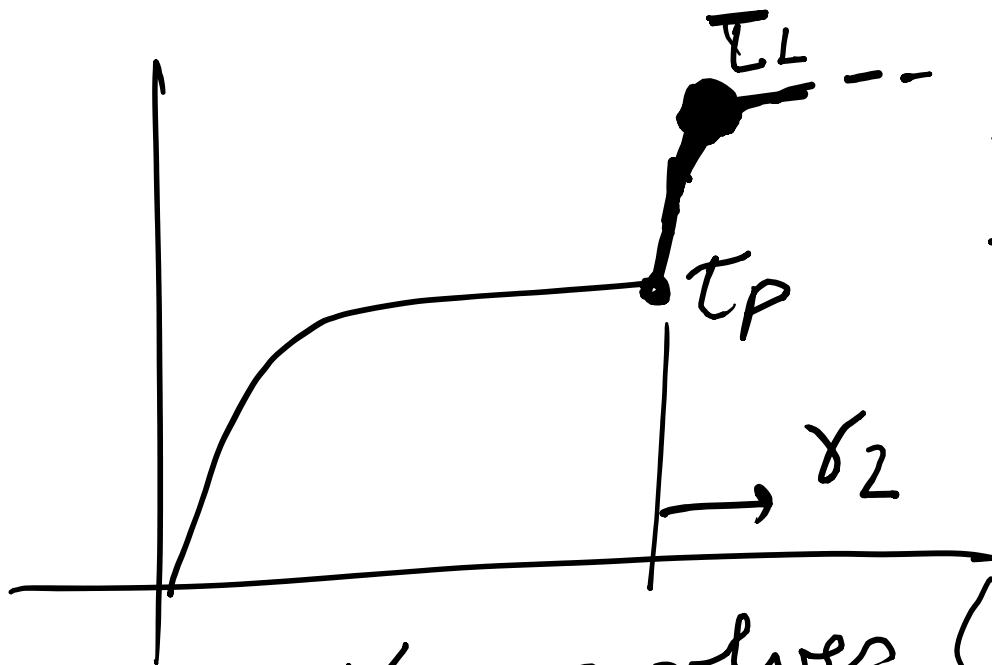
$\underbrace{\hspace{10em}}$   
from step I  
loading

$$\tau_1 = \underbrace{h_{11} \gamma_1^{\max}}_{\text{initial strength from step I}} + \underbrace{h_{12} \gamma_2}_{\text{latent hardening}}$$

During step II loading,

$$L_{ij} = \frac{T_L}{T_p} = \frac{h_{22} \delta_2 + h_{21} \delta_1^{\max}}{h_{11} \delta_1^{\max} + h_{12} \delta_2}$$

When  $\delta_2 = 0$  at beginning Step II  
one can evaluate the  
 $h_{21} / h_{11}$  ratio.

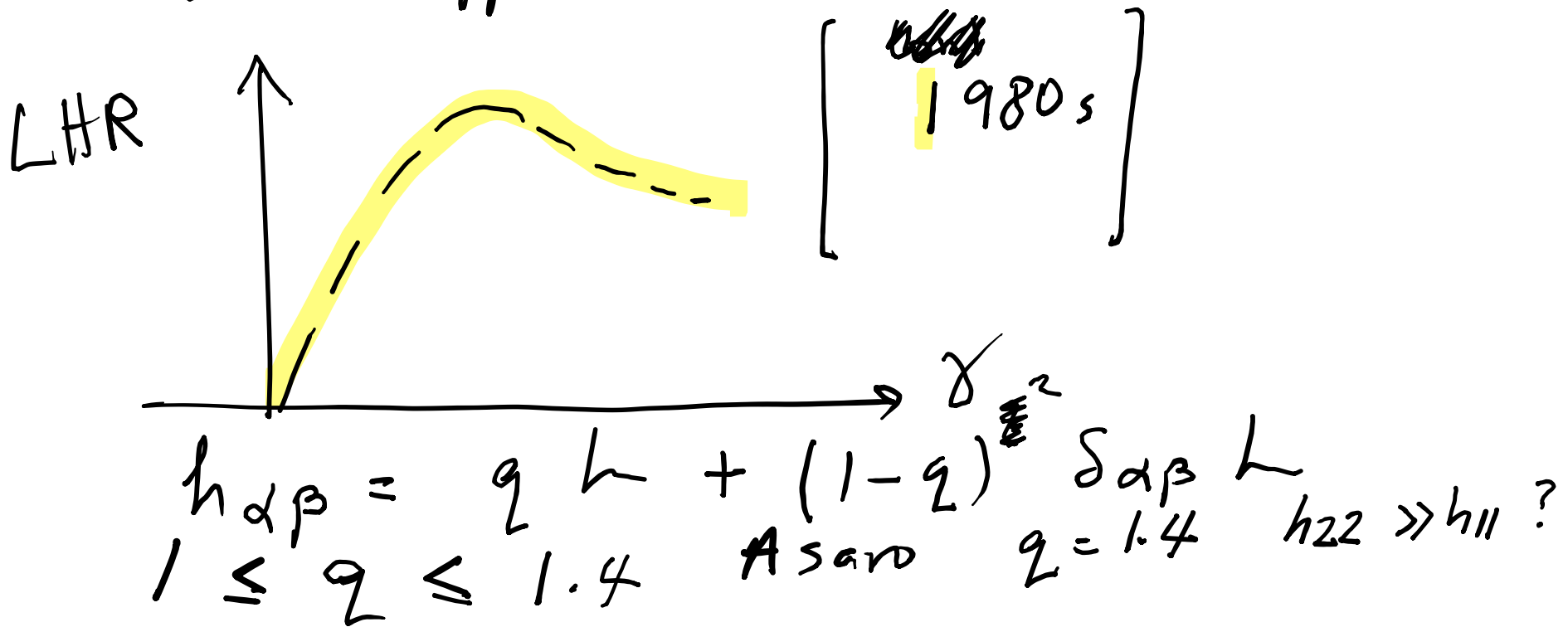


$$\frac{T_L}{T_p} = \text{Known for } \gamma_2 = 0$$

As  $\gamma_2$  evolves (if  $h_{12}$  is assumed same as  $h_{21}$ ) one can use this equation to solve for  $h_{22}$ .

In fact  $LHR$  increases first reaches a peak. Exp. results

imply that  $h_{22}$  may not be const. it decreases with applied strain.



Saneer, WARD 3D CP option Monday

- 1) ✓ Latent hard. ✓
- 2) glom.  $k_0 = 0$
- 3) Stress rates (  $\sigma =$  )



































































































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